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## On a regular $\psi$ -fractional Sturm-Liouville problem

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**Abstract:** In this short paper, we consider a  $\psi$ -fractional Sturm-Liouville eigenvalue problem by using left  $\psi$ -Caputo and right  $\psi$ -Riemann-Liouville fractional derivatives. We study the main properties of the eigenfunctions and the eigenvalues of the associated fractional boundary problem.

**keywords:**  $\psi$ -fractional derivatives;  $\psi$ -fractional Sturm-Liouville problem; Eigenvalues; Eigenfunctions.

**MSC2010:** 34B24; 26A33; 34L10.

## 1 Introduction

The classical regular Sturm-Liouville (S-L) problem is associated with the real second-order linear ordinary differential equation of the form:

$$\frac{d}{dx} \left( p(x) \frac{dy}{dx} \right) + q(x) y(x) = \lambda w(x) y(x),$$

where  $p(x), w(x) > 0$ , and  $p(x), p'(x), q(x), w(x)$  are continuous functions on  $(a, b)$ . The function  $w(x)$ , sometimes denoted by  $r(x)$ , is called the weight or density function. The unknown function  $y(x)$  is continuous and differentiable on  $(a, b)$ . In addition,  $y(x)$  is required to satisfy some boundary conditions of the form  $c_1 y(a) + c_2 y'(a) = 0$  and  $d_1 y(b) + d_2 y'(b) = 0$ , with  $c_1^2 + c_2^2 \neq 0$  and  $d_1^2 + d_2^2 \neq 0$ . Finding the values of  $\lambda$ , known as eigenvalues, for which there exists a non-trivial solution  $y(x)$ , called eigenfunction associated to  $\lambda$ , constitutes the classical S-L problem. Fractional versions of this problem were already considered in the literature (see e.g. [3, 4] and references therein). Motivated by the unification given by the  $\psi$ -fractional calculus we study a regular  $\psi$ -fractional S-L problem.

## 2 Preliminaries

In this section, we introduce the concepts related to the  $\psi$ -fractional calculus necessary to this work (for more details see [5] and references therein).

**Definition 1** (cf. [5]) Let  $(a, b)$  be a finite or infinite interval on the real line  $\mathbb{R}$  and  $\alpha > 0$ . Also let  $\psi$  be a monotone increasing and positive function on  $(a, b)$ , having a continuous derivative  $\psi'$  in  $(a, b)$ . The left- and right-sided Riemann-Liouville (RL) fractional integrals of a function  $f$  with respect to another function  $\psi$  on  $(a, b)$  are given respectively by

$$I_{a^+}^{\alpha;\psi} f(x) = \frac{1}{\Gamma(\alpha)} \int_a^x \psi'(t) (\psi(x) - \psi(t))^{\alpha-1} f(t) dt, \quad x > a, \quad (1)$$

$$I_{b^-}^{\alpha;\psi} f(x) = \frac{1}{\Gamma(\alpha)} \int_x^b \psi'(t) (\psi(t) - \psi(x))^{\alpha-1} f(t) dt, \quad x < b. \quad (2)$$

Assuming further that  $f, \psi \in C^n(a, b)$ , where  $n = [\alpha] + 1$ , and  $\psi'(x) \neq 0$ , for all  $x \in (a, b)$ , the corresponding inverse operators, i.e., the left- and right-sided  $\psi$ -RL fractional derivatives of order  $\alpha > 0$ , are defined respectively by

$$D_{a^+}^{\alpha;\psi} f(x) = \left( \frac{1}{\psi'(x)} \frac{d}{dx} \right)^n I_{a^+}^{n-\alpha;\psi} f(x) \quad \text{and} \quad D_{b^-}^{\alpha;\psi} f(x) = \left( -\frac{1}{\psi'(x)} \frac{d}{dx} \right)^n I_{b^-}^{n-\alpha;\psi} f(x). \quad (3)$$

Moreover, the left- and right-sided  $\psi$ -Caputo fractional derivatives of order  $\alpha > 0$ , are defined respectively by

$${}^C D_{a^+}^{\alpha;\psi} f(x) = I_{a^+}^{n-\alpha;\psi} \left( \frac{1}{\psi'(x)} \frac{d}{dx} \right)^n f(x) \quad \text{and} \quad {}^C D_{b^-}^{\alpha;\psi} f(x) = I_{b^-}^{n-\alpha;\psi} \left( -\frac{1}{\psi'(x)} \frac{d}{dx} \right)^n f(x). \quad (4)$$

Using the concepts of  $\psi$ -RL and  $\psi$ -Caputo fractional derivatives we can encompass in one definition several fractional derivatives known in the literature. In [5], the authors listed several fractional derivatives that can be obtained for specific choices of the function  $\psi$  such as Caputo, RL, Hadamard, Katugampola, Chen, Jumarie, Prabhakar, Erdélyi-Kober, Weyl, among others.

**Proposition 1** Let  $\alpha, \beta > 0$  and  $f \in C([a, b])$ . Then the following relations hold

$$\begin{aligned} I_{a^+}^{\alpha;\psi} I_{a^+}^{\beta;\psi} &= I_{a^+}^{\alpha+\beta;\psi}, & I_{b^-}^{\alpha;\psi} I_{b^-}^{\beta;\psi} &= I_{b^-}^{\alpha+\beta;\psi}, \\ D_{a^+}^{\alpha;\psi} I_{a^+}^{\alpha;\psi} f(x) &= f(x), & D_{b^-}^{\alpha;\psi} I_{b^-}^{\alpha;\psi} f(x) &= f(x), \\ {}^C D_{a^+}^{\alpha;\psi} I_{a^+}^{\alpha;\psi} f(x) &= f(x), & {}^C D_{b^-}^{\alpha;\psi} I_{b^-}^{\alpha;\psi} f(x) &= f(x). \end{aligned}$$

Moreover, for  $f \in C^n([a, b])$ ,  $n = [\alpha] + 1$ , and putting  $f_\psi^{[k]}(x) = \left( \frac{1}{\psi'(x)} \frac{d}{dx} \right)^k f(x)$ , we have

$$\begin{aligned} {}^C D_{a^+}^{\alpha;\psi} f(x) &= D_{a^+}^{\alpha;\psi} \left[ f(x) - \sum_{k=0}^{n-1} \frac{f_\psi^{[k]}(a)}{k!} (\psi(x) - \psi(a))^k \right], \\ I_{a^+}^{\alpha;\psi} {}^C D_{a^+}^{\alpha;\psi} f(x) &= f(x) - \sum_{k=0}^{n-1} \frac{f_\psi^{[k]}(a)}{k!} (\psi(x) - \psi(a))^k. \end{aligned}$$

**Proposition 2** For  $\alpha \in (0, 1)$ , the fractional differential operators (3) and (4) satisfy the following  $\psi$ -fractional integration by parts:

$$\int_a^b f(x) D_{a^+}^{\alpha;\psi} g(x) \psi'(x) dx = \int_a^b {}^C D_{b^-}^{\alpha;\psi} f(x) g(x) \psi'(x) dx + f(x) I_{a^+}^{1-\alpha;\psi} g(x) \Big|_{x=a}^{x=b}, \quad (5)$$

$$\int_a^b f(x) D_{b^-}^{\alpha;\psi} g(x) \psi'(x) dx = \int_a^b {}^C D_{a^+}^{\alpha;\psi} f(x) g(x) \psi'(x) dx - f(x) I_{b^-}^{1-\alpha;\psi} g(x) \Big|_{x=a}^{x=b}. \quad (6)$$

For more properties of these  $\psi$ -fractional operators see [5] and references therein.

### 3 Regular $\psi$ -fractional Sturm-Liouville problem

We consider the following  $\psi$ -fractional S-L equation

$$-D_{b^-}^{\alpha;\psi} \left( p(x) {}^C D_{a^+}^{\alpha;\psi} f(x) \right) + q(x)f(x) = \lambda w(x) f(x) \quad (7)$$

subject to the following boundary conditions

$$c_1 f(a) + c_2 I_{b^-}^{1-\alpha;\psi} \left( p(x) {}^C D_{a^+}^{\alpha;\psi} f(x) \right) \Big|_{x=a} = 0, \quad c_1^2 + c_2^2 \neq 0, \quad (8)$$

$$d_1 f(b) + d_2 I_{b^-}^{1-\alpha;\psi} \left( p(x) {}^C D_{a^+}^{\alpha;\psi} f(x) \right) \Big|_{x=b} = 0, \quad d_1^2 + d_2^2 \neq 0, \quad (9)$$

where  $\alpha \in (0, 1)$ ,  $x \in [a, b]$ ,  $p(x), w(x) > 0$ , for all  $x \in [a, b]$ , and  $p, q, w$  are real-valued continuous functions in  $[a, b]$ . Our problem consists in finding the values of  $\lambda$  such that the boundary-value problem (7)-(9) has a non-trivial solution. We denote by  $\mathcal{L}^{\alpha;\psi}$  the  $\psi$ -fractional S-L operator associated to problem (7) given by

$$\mathcal{L}^{\alpha;\psi} := -D_{b^-}^{\alpha;\psi} \left( p(x) {}^C D_{a^+}^{\alpha;\psi} \right) + q(x).$$

As suggested by the  $\psi$ -fractional integration by parts of Proposition 2 we will work in the weighted space  $L_1([a, b], \psi'(x)dx)$ . Using (6) and the boundary conditions (8)-(9) one can prove the following theorem.

**Theorem 3** *The eigenvalues of our  $\psi$ -fractional problem (7)-(9) are real. Moreover, the eigenfunctions corresponding to distinct eigenvalues are orthogonal with respect to the sensity function  $w(x)$  on  $[a, b]$ , that is*

$$\int_a^b f_{\lambda_1}(x) f_{\lambda_2}(x) w(x) \psi'(x) dx, \quad \lambda_1 \neq \lambda_2.$$

Our  $\psi$ -fractional S-L equation with boundary conditions can be written in integral form as the next proposition shows.

**Proposition 4** *Let  $\alpha > 1/2$  and  $f_\lambda$  be an eigenfunction associated to the eigenvalue  $\lambda$ . On the space  $C([a, b])$ , the  $\psi$ -fractional S-L problem is equivalent to the integral equation*

$$f_\lambda(x) = -I_{a^+}^{\alpha;\psi} \left( \frac{1}{p(x)} I_{b^-}^{\alpha;\psi} F_\lambda(f) \right) + A(x) \int_a^b F_\lambda(f) dx + B(x) \left( I_{a^+}^{\alpha;\psi} \left( \frac{1}{p(x)} I_{b^-}^{\alpha;\psi} F_\lambda(f) \right) \right) \Big|_{x=b} \quad (10)$$

where the coefficients  $A(x)$  and  $B(x)$  are given by

$$A(x) = \frac{c_2}{\Delta} \left( d_2 + d_1 \left( Y - I_{a^+}^{\alpha;\psi} \left( \frac{(\psi(b) - \psi(x))^{\alpha-1}}{p(x) \Gamma(\alpha)} \right) \right) \right),$$

$$B(x) = \frac{d_1}{\Delta} \left( c_1 I_{a^+}^{\alpha;\psi} \left( \frac{(\psi(b) - \psi(x))^{\alpha-1}}{p(x) \Gamma(\alpha)} \right) - c_2 \right),$$

with  $\Delta = c_1(d_1 Y + d_2) - c_2 d_1 \neq 0$ ,  $Y = I_{a^+}^{\alpha;\psi} \left( \frac{(\psi(b) - \psi(x))^{\alpha-1}}{p(x) \Gamma(\alpha)} \right) \Big|_{x=b}$ , and  $F_\lambda(f) = q(x)f_\lambda(x) - \lambda w(x)f_\lambda(x)$ .

**Sketch of the proof:** Using composition rules we can write the equation (7) as follows:

$$D_{b^-}^{\alpha;\psi} \left( p(x) {}^C D_{a^+}^{\alpha;\psi} \left[ f_\lambda(x) + I_{a^+}^{\alpha;\psi} \left( \frac{1}{p(x)} I_{b^-}^{\alpha;\psi} F_\lambda(f) \right) \right] \right) = 0. \quad (11)$$

On the space  $C([a, b])$  a general solution of (11) is given by

$$f_\lambda(x) + I_{a^+}^{\alpha;\psi} \left( \frac{1}{p(x)} I_{b^-}^{\alpha;\psi} F_\lambda(f) \right) = \xi_1 + \xi_2 I_{a^+}^{\alpha;\psi} \left( \frac{(\psi(b) - \psi(x))^{\alpha-1}}{p(x) \Gamma(\alpha)} \right), \quad \xi_1, \xi_2 \in \mathbb{R}. \quad (12)$$

From [3, Eq. (34)] the fractional integral that is multiplied by  $\xi_2$  gives a continuous function on  $[a, b]$  only for  $1/2 < \alpha < 1$ . From (12) and using composition rules we obtain:

$$I_{b^-}^{1-\alpha;\psi} \left( p(x) {}^C D_{a^+}^{\alpha;\psi} f_\lambda(x) \right) = \xi_2 - I_{b^-}^{1;\psi} F_\lambda(f). \quad (13)$$

Finally, from (13) and (12) and using the boundary conditions (8)-(9) we can relate the coefficients  $\xi_1, \xi_2$  with the values  $c_j, d_j, j = 1, 2$  in the boundary conditions to obtain (10).

Our final theorem gives the conditions under which the eigenfunctions exist and are unique. The proof follows similar arguments as the proof of [3, Thm. 9].

**Theorem 5** *Let  $\alpha > 1/2$  and assume that  $\Delta \neq 0$ . Then our  $\psi$ -fractional problem has unique continuous eigenfunctions  $f_\lambda$  if the fixed point condition holds*

$$\|q(x) + \lambda\omega(x)\| < \frac{m_p}{\|\varphi(x)\| + \|B(x)\|\varphi(b) + m_p\|A(x)\|(\psi(b) - \psi(a))},$$

where  $\varphi(x) = I_{a^+}^{\alpha;\psi} \left( \frac{(\psi(b) - \psi(x))^{\alpha-1}}{p(x) \Gamma(\alpha)} \right)$  and  $m_p = \min_{x \in [a, b]} |p(x)|$ .

## 4 Conclusion

In this work we showed how to use the  $\psi$ -fractional calculus to study a  $\psi$ -fractional S-L problem.

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